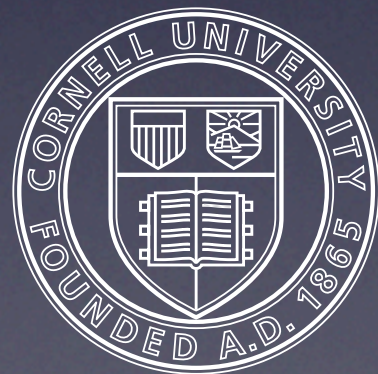


The flavor of the holographic pGB

Andreas Weiler
Brookhaven Forum 2008



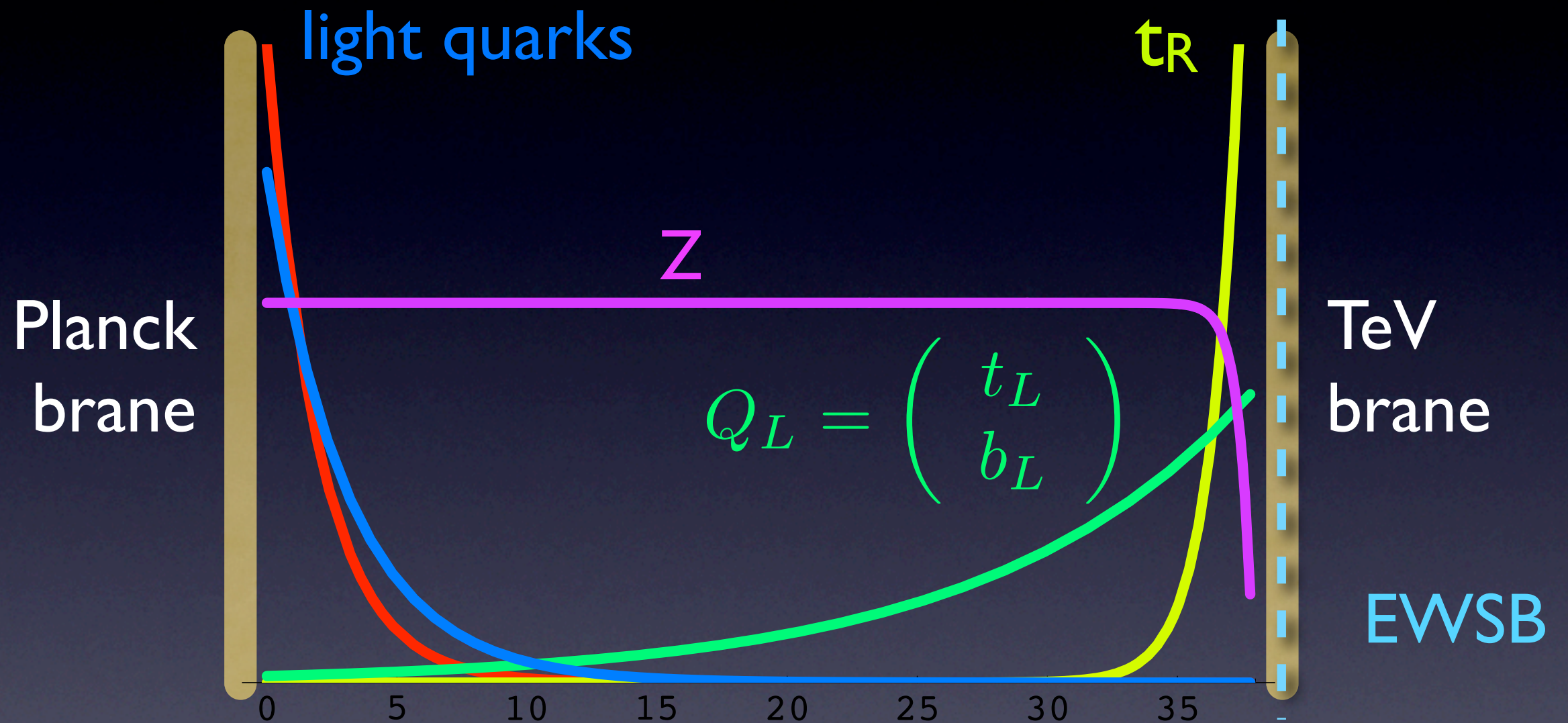
Cornell University

based on papers with:

C. Csaki and A. Falkowski (pGB flavor, $U(1)$'s);
C. Csaki, Y. Grossman, G. Perez, and Z. Surujon

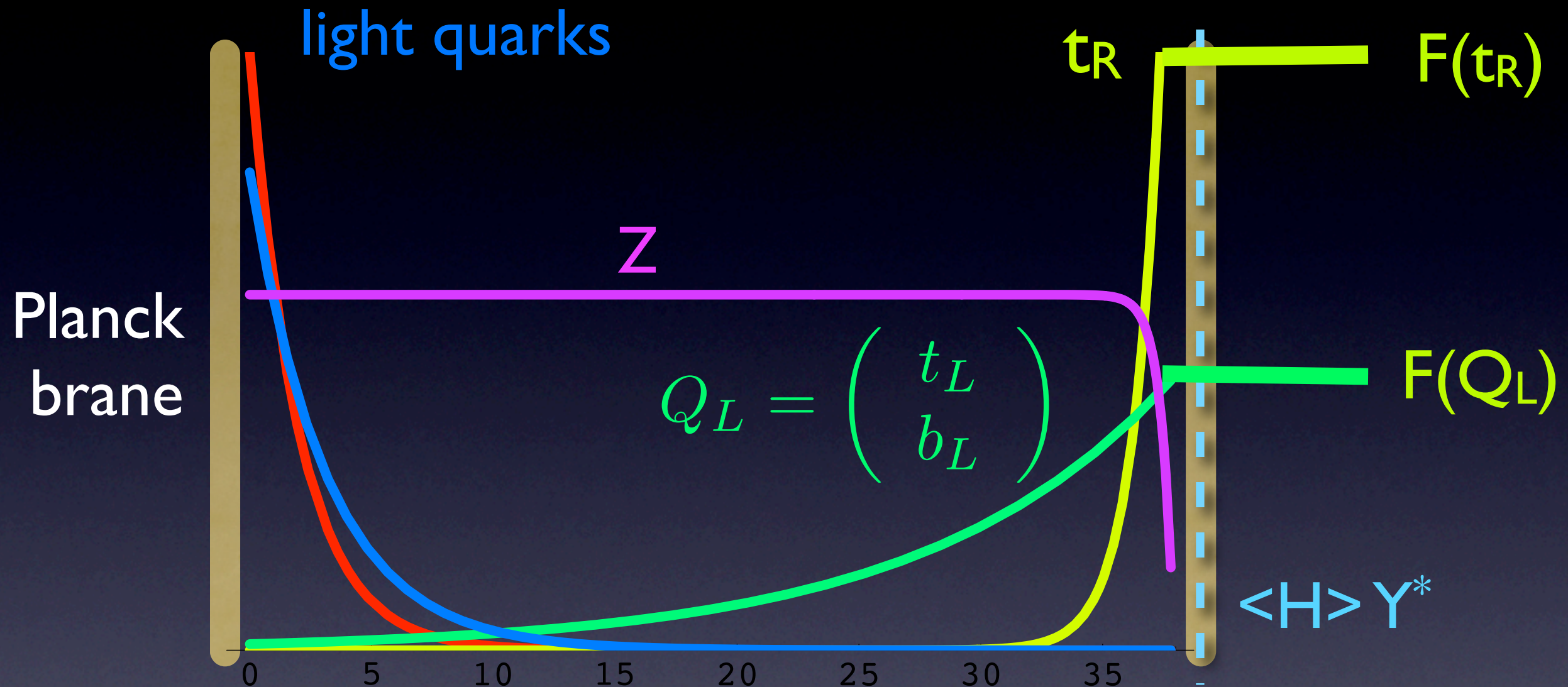
Hierarchies without symmetries

Arkani-Hamed, Schmaltz; Grossman, Neubert;
Gherghetta, Pomarol



warped extra dimensions as a theory of flavor

Wavefunction overlap generates hierarchies



F = wave function @ IR brane :

$$F \sim (\text{TeV/Planck})^{2c-1} \quad c > 1/2, \quad c = \text{bulk mass}$$

$$F \sim \sqrt{1-2c} \quad c < 1/2$$

RS has a flavor problem

LR chiral contributions to CPV in K - \bar{K} mixing
generically require $m_{KK} > 11 \text{ TeV}^+$

But there is also fine-tuning in EWSB...

+ terms and conditions apply

(scalar bulk Higgs with less perturbative control $m > 5 \text{ TeV}$ possible)

RS and little hierarchy problem

Precision electroweak data suggests

- o light Higgs ($m_H < 200 \text{ GeV}$)
- o (S,T,U,...) new contributions $\Lambda > 5 \text{ TeV}$

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Large UV sensitive contribution to Higgs mass.

Top loops e.g. induce

$$m_H^2 \sim -3 \lambda_t^2 \frac{1}{8\pi^2} \Lambda^2$$

Significant fine-tuning if not taken care of:

RS with Higgs on the brane or scalar bulk Higgs suffers from little hierarchy problem!

Solution: pGB Higgs models

Agashe, Contino, Pomarol 05

Simple model with

- o custodial symmetry
- o A_5 zero mode $\in SO(5)/SO(4) = \text{Higgs}$
- o small corrections to S, T, U, Z_{bb}



Solution: pGB Higgs models

Agashe, Contino, Pomarol 05

Simple model with

- o custodial symmetry
- o A_5 zero mode $\in SO(5)/SO(4) = \text{Higgs}$
- o s

Dual to pGB composite
Higgs (Georgi, Kaplan '83)

R_{ir}

Planck
brane

AdS₅
 $SO(5) \times U(1)_x$

TeV
brane

$SU(2) \times U(1)_Y$

$SO(4) \times U(1)_x$

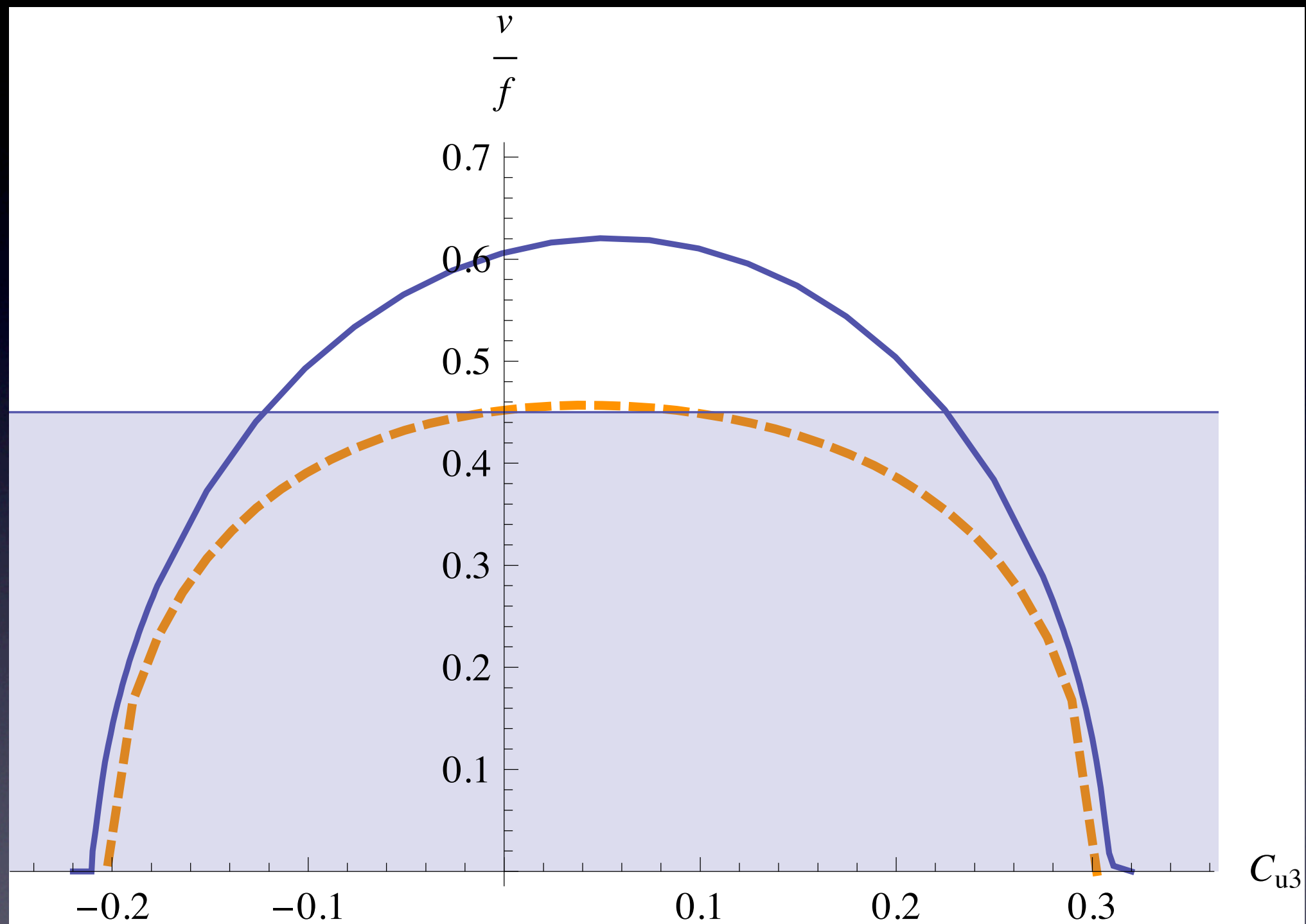
Calculable radiative EWSB

$A_5 = \text{Higgs}$: Non-local Coleman-Weinberg
induces potential and $\langle A_5 \rangle \neq 0$

$$V(v) = \frac{3}{32\pi^2} \int_0^\infty dt t [-4 \log \rho_t(-t) + 2 \log \rho_W(-t) + \log \rho_Z(-t)]$$

UV finite, depends on **5D fermion mass sector**

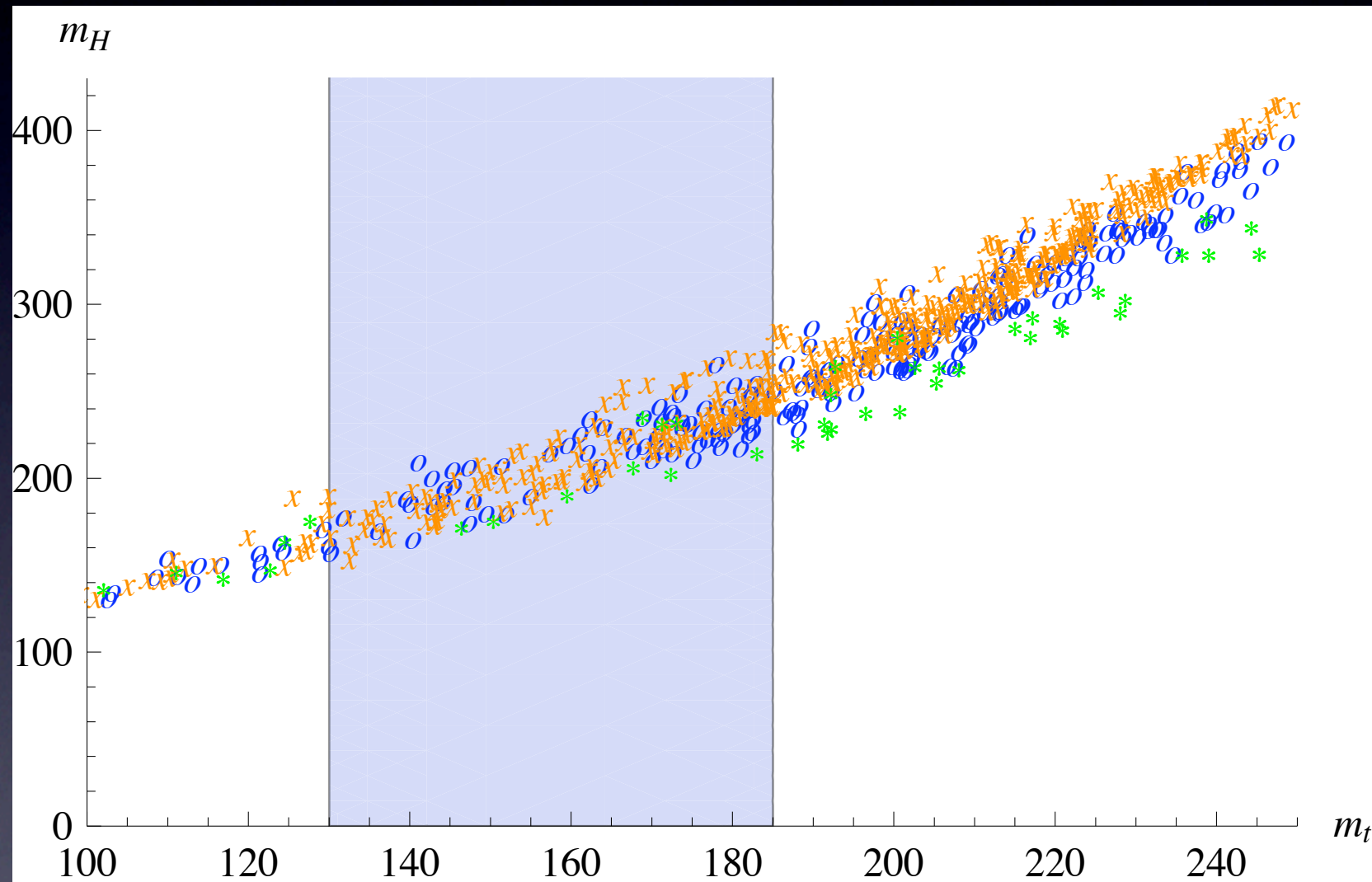
$\langle H \rangle$ / (new physics scale)



c_{u3} : bulk mass of top

Realistic EWSB only for correlated parameter set

m_H



m_t

Fermion Masses in Gauge-Higgs models

A_5 = Higgs being a gauge field couples only to fields in the same multiplet.

Add boundary mixing terms:
Zero modes lives in multiple representations
→ kinetic mixing.

Yukawa = bulk gauge coupling g_*

Mass terms from gauge interactions

Csaki, Falkowski, AW

Some freedom to embed fermion content

Example here: 4 (spinor) of SO(5)

$$\Psi_q = \begin{pmatrix} q_q[+, +] \\ u_q^c[-, +] \\ d_q^c[-, +] \end{pmatrix} \quad \Psi_u = \begin{pmatrix} q_u[+, -] \\ u_u^c[-, -] \\ d_u^c[+, -] \end{pmatrix} \quad \Psi_d = \begin{pmatrix} q_d[+, -] \\ u_d^c[+, -] \\ d_d^c[-, -] \end{pmatrix}$$

I) = chiral zero modes

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$$\begin{array}{c} \text{A}_5 \end{array} \Rightarrow \begin{pmatrix} q_q[+, +] \\ u_q^c[-, +] \\ d_q^c[-, +] \end{pmatrix} \quad \Psi_u = \begin{pmatrix} q_u[+, -] \\ u_u^c[-, -] \\ d_u^c[+, -] \end{pmatrix} \quad \Psi_d = \begin{pmatrix} q_d[+, -] \\ u_d^c[+, -] \\ d_d^c[-, -] \end{pmatrix}$$

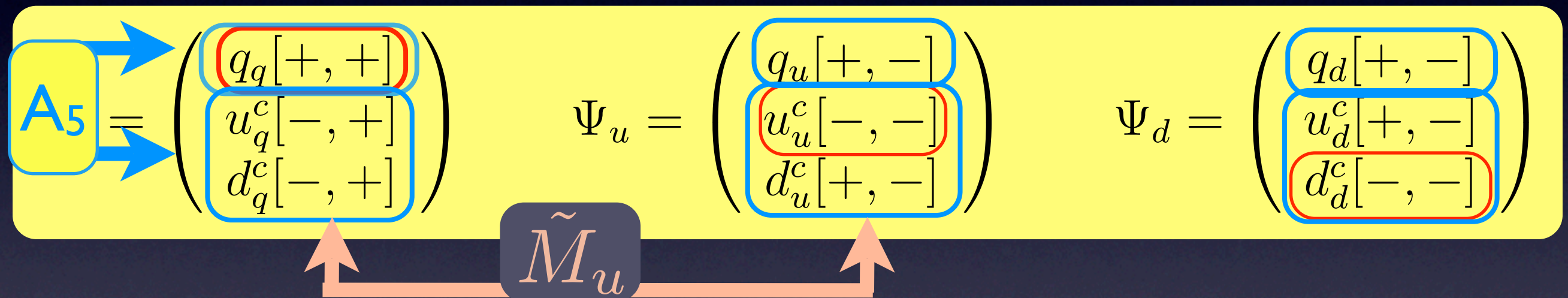
- 1) = chiral zero modes
- 2) $\langle \text{A}_5 \rangle$ marries fields in same multiplet

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- 1) = chiral zero modes
- 2) $\langle A_5 \rangle$ marries fields in same multiplet
- 3) SO(4) invariant **brane mixings** mix multiplets

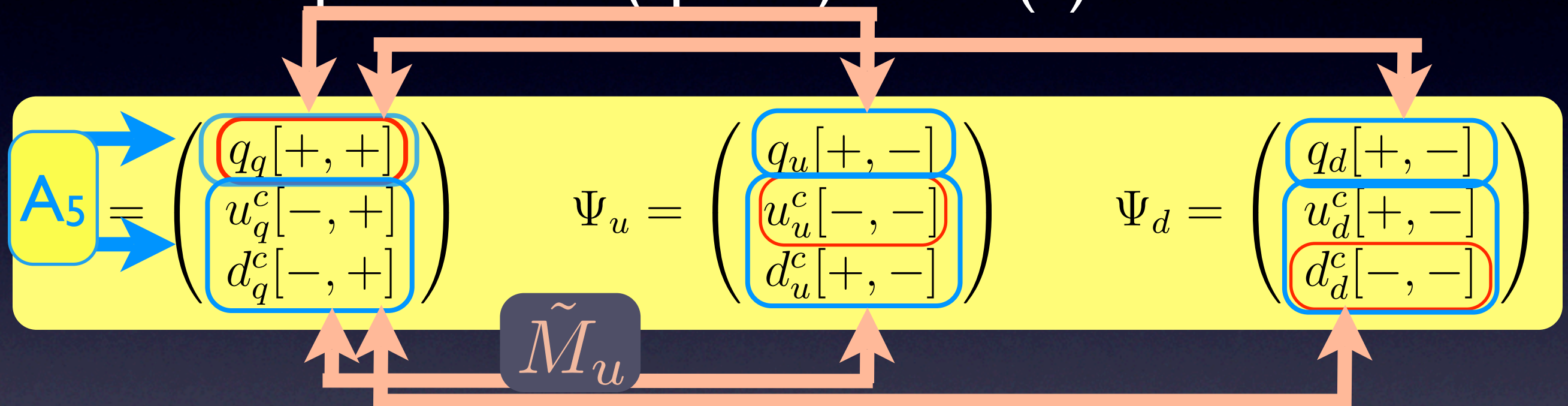
$$\mathcal{L}_{IR} = - \left(\frac{R}{R'} \right)^4 \left[\tilde{m}_u \chi_{q_q} \psi_{q_u} + \tilde{m}_d \chi_{q_q} \psi_{q_d} + \tilde{M}_u (\chi_{u_q^c} \psi_{u_u^c} + \chi_{d_q^c} \psi_{d_u^c}) + \tilde{M}_d (\chi_{u_q^c} \psi_{u_d^c} + \chi_{d_q^c} \psi_{d_d^c}) \right]$$

Mass terms from gauge interactions

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Effective mass terms pGB model

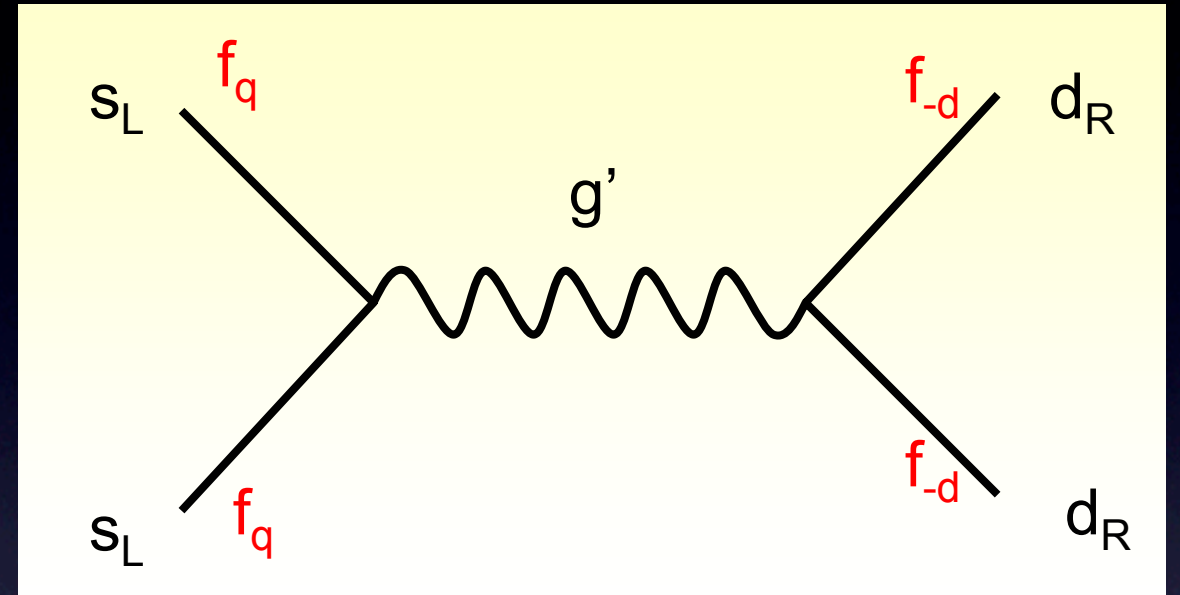
$$m_u^{SM} = \frac{g_* v}{2\sqrt{2}} H_q f_q (\tilde{m}_u - \tilde{M}_u) f_{-u} H_u$$

$$m_d^{SM} = \frac{g_* v}{2\sqrt{2}} H_q f_q (\tilde{m}_d - \tilde{M}_d) f_{-d} H_d$$

Effective 4 fermi operators

Csaki, Falkowski, AW

Integrating out the
KK gluon



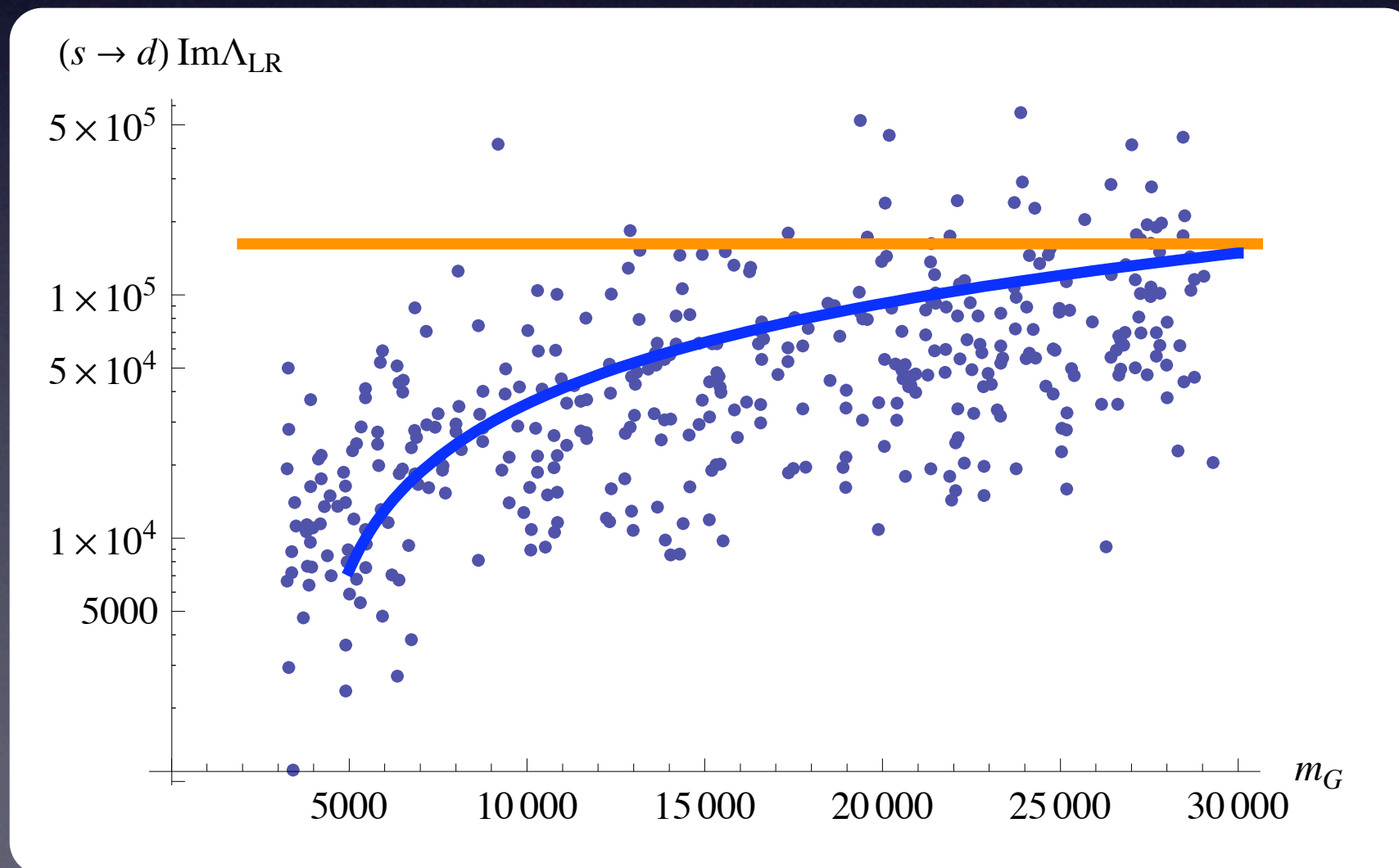
$$\begin{aligned}\mathcal{H} &= \frac{1}{M_G^2} \left[\frac{1}{6} g_L^{ij} g_L^{kl} (\bar{q}_L^{i\alpha} \gamma_\mu q_{L\alpha}^j) (\bar{q}_L^{k\beta} \gamma^\mu q_{L\beta}^l) - g_R^{ij} g_L^{kl} \left((\bar{q}_R^{i\alpha} q_{L\alpha}^k) (\bar{q}_L^{l\beta} q_{R\beta}^j) - \frac{1}{3} (\bar{q}_R^{i\alpha} q_{L\beta}^l) (\bar{q}_L^{k\beta} q_{R\alpha}^j) \right) \right] \\ &= C^1(M_G) (\bar{q}_L^{i\alpha} \gamma_\mu q_{L\alpha}^j) (\bar{q}_L^{k\beta} \gamma^\mu q_{L\beta}^l) + C^4(M_G) (\bar{q}_R^{i\alpha} q_{L\alpha}^k) (\bar{q}_L^{l\beta} q_{R\beta}^j) + C^5(M_G) (\bar{q}_R^{i\alpha} q_{L\beta}^l) (\bar{q}_L^{k\beta} q_{R\alpha}^j)\end{aligned}$$

$$C_{4K}^{RS} \sim \frac{g_{s*}^2}{M_G^2} f_{q_1} f_{q_2} f_{-d_1} f_{-d_2} \sim \frac{1}{M_G^2} \frac{g_{s*}^2}{Y_*^2} \frac{2m_d m_s}{v^2}$$

Four fermi operators

$$C_K^4 \sim \frac{1}{M_G^2} \frac{g_{s*}^2}{g_*^2} \frac{8m_d m_s}{v^2} \frac{1 + m^2}{\tilde{m}_d^2}$$

pGB worse than RS: $Y^* \leftrightarrow g^* / 2$, $M_{KK} > 30 \text{ TeV}$



Low KK scale w/o adding flavor structure

+ live with fine-tuned Yukawas (large radiative corrections)

Blanke, Buras, Dulling, Gori, AW; Casagrande, Goertz, Haisch, Neubert, Pfoh

or

Agashe, Azatov, Zhu

+ bulk Higgs model (**not applicable to pGB**),
push Yukawa to perturbative limit $Y^* > 6$ and
 g_{s^*} as small as possible (1-loop matching)

$$M_{KK} > \frac{g_{s^*}}{Y^*} \frac{\sqrt{2m_d m_s}}{v} \Lambda_4$$

With some tuning **$M_{KK} \sim 5 \text{ TeV}$** possible
Testable at LHC? Little hierarchy?

Low **KK** scale by adding flavor structure

Csaki, Falkowski, AW

- + Propose $U(1)_d \times U(1)_q$ for quark representation with custodial protection of Z_{bb} of pGB
- Key ingredient: two rep.'s for (q_u, q_d) for Q_L

Low KK scale by adding flavor structure

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BC on UV:

$$\theta q_u - q_d = 0$$

Forces θ to be diagonal

	$U(1)_d$ (-, +)	$U(1)_q$ (+, -)
q_u	0	$\begin{pmatrix} q_1 \\ q_2 \\ q_3 \end{pmatrix}$
u	0	0
q_d	$\begin{pmatrix} \tilde{q}_1 \\ \tilde{q}_2 \\ \tilde{q}_3 \end{pmatrix}$	$\begin{pmatrix} q_1 \\ q_2 \\ q_3 \end{pmatrix}$
d	$\begin{pmatrix} \tilde{q}_1 \\ \tilde{q}_2 \\ \tilde{q}_3 \end{pmatrix}$	0

Low KK scale by adding flavor structure II

Csaki, Falkowski, AW

+ Propose $U(1)_d \times U(1)_q$ for quark representation with custodial protection of Z_{bb} of pGB

Key ingredient: two rep.'s for (q_u, q_d) for Q_L

IR masses:

$$\frac{v}{\sqrt{2}} f_{q_u} Y_u f_{-u}$$

$$\frac{v}{\sqrt{2}} f_{q_d} Y_d f_{-d}$$

Y_d diagonal

	$U(1)_d$ (-, +)	$U(1)_q$ (+, -)
q_u	0	$\begin{pmatrix} q_1 \\ q_2 \\ q_3 \end{pmatrix}$
u	0	0
q_d	$\begin{pmatrix} \tilde{q}_1 \\ \tilde{q}_2 \\ \tilde{q}_3 \end{pmatrix}$	$\begin{pmatrix} q_1 \\ q_2 \\ q_3 \end{pmatrix}$
d	$\begin{pmatrix} \tilde{q}_1 \\ \tilde{q}_2 \\ \tilde{q}_3 \end{pmatrix}$	0

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- + Propose $U(1)_d \times U(1)_q$ for quark representation with custodial protection of Z_{bb} of pGB
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All flavor violation in up-sector, constraint from D - \underline{D} bar mixing: $M_{KK} > 1 \text{ TeV}$

$U(1)_d \times U(1)_q$ gauge bosons give additional contributions, need them to be almost global
 $g_5^D < 1/50 g_{QCD}$.

Conclusions

- RS GIM suppresses most of the dangerous FCNCs
- Contributions to ε_K with LR chirality typically too large
- Bulk Higgs $m_{KK} > 5 \text{ TeV}$ (best cases)
 Brane Higgs $m_{KK} > 11 \text{ TeV}$
 pGB Higgs $m_{KK} > 15 \text{ TeV}$
- Additional mechanisms needed, e.g.
 horizontal $U(1)$'s to allow $m_{KK} \sim 1\text{-}2 \text{ TeV}$

Low KK scale by adding flavor structure

Cacciapaglia, Csaki, Galloway, Marandella, Terning, AV

+ exact GIM structure

flavor symmetry in bulk and IR brane, UV kinetic terms generate flavor, no explanation for fermion masses (likely the only way for Higgsless)

Santiago

+ Minimal flavor protection bulk $U(3)$ flavor symmetry in d_R sector (radiatively unstable)

Fitzpatrick, Randall, Perez

+ 5D MFV only two flavor spurions (Y_U, Y_D) Need to align bulk and brane matrices by hand. Can we really avoid tuning?